

# **Lectures on the Physics of the Nucleus**

**FZ Dresden-Rossendorf, Feb-April 2008**

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## **Overview:**

- **show possible subjects of lecture**
- **find out and select, what is of interest to the audience**

## A. Structure

### 1. Single particle models

bulk properties, nuclear matter, Fermi gas  
indep. part. shell model (SM), spin-orbit, SM states, s.p. transitions  
deformed SM, Nilsson model  
deformation energy surfaces, Strutinsky shell corrections  
pairing, quasi-particles, BCS

### 2. Compound nuclear model

level densities, mean free path

### 3. NN-interaction

2N-states, bound, scattering  
realistic interactions, meson exchange, chiral interactions  
effective interactions, T-, G-matrix,  $V_{lowk}$

### 4. Many-body methods

SM with residual interactions, no-core SM, Monte-Carlo SM, Greenfct.-MC  
Hartree-Fock, HF-Bogoliubov, Brueckner-HF  
field theoretical approach, Quantumhydrodynamics (QHD), Rel. Mean field (RMF)

### 5. Decay modes

el.magn. decay, multipole-moments and -operators, transition rates  
 $\alpha$ -decay  
 $\beta$ -decay, weak interaction, parity, Fermi-, Gamow-Teller transitions  
fission

### 6. Collective motion: Rotations

dynamic and static deformations, rotations, vibrations, coupling, bands  
(microscopic) HF, HFB, constrained HBF  
cranking model, moment of inertia,  
high spin, back bending

### 7. Collective motion: vibrations

harmonic vibrations, classification of modes, giant resonances (GR), exotic modes  
(micr.) TDA, RPA, TDHF, GCM  
strength functions, sum rules  
boson expansions, Interact. Boson approximation (IBA)

## B. Reactions

### 1. Reaction theory

scattering states, wave packet scattering, cross sections classical and qm,  
Born approx. (BA), partial waves, scattering amplitude  
scattering lengths, effective range theory  
reactions, two-potential formula, DWBA  
operator formalism, S-matrix  
R-matrix approach

### 2. Compound nucleus scattering

resonances, overlapping resonances, fluctuations (Ericson),  
CN scattering, Hauser-Feshbach formalism  
average scattering, optical model

### 3. Optical Model

Projektion formalism (Feshbach)  
complex potential, mean free path, dispersion relations  
parametrizations, folding,

### 4. Direct reactions

DWBA, angular distributions, polarization  
inelastic scattering, collective strength functions  
charge exchange,  $\beta$ -strength functions  
transfer, (1 and many nucleon), spectroscopic factors,  
multistep reactions, CC, CCBA, CRC

### 5. Statistical reactions

pre-equilibrium reactions  
exciton models  
quantummech. pre-eq. theories (FKK, TUL)

### 6. Heavy ion collisions

semiclassical features  
deep inelastic scattering, dissipation  
transport approach, BUU equations  
fragmentation

# 0. Introductory Remarks

Modell des Kerns: stark wechselwirkendes Vielteilchensystem

Vielteilchenaspekt  
 N-Teilchensystem:  $1 \leq N \lesssim$  einige Hundert  
 Hamiltonfunktion  

$$H = \sum_{i=1}^N T_i + \sum_{i,j=1}^N V_{ij}$$
 ↑  
 mesoskopisches System  
 Spezifiziert, lösbar?

steigende Energieskala ↓

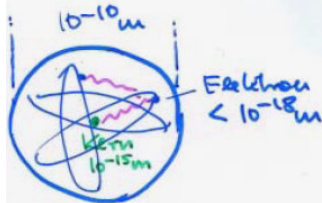
1. Nukleonen, $V_{NN}$ („klass.“ Kernphysik)	
2. Nukleonen, Mesonen (effekt. Feldtheorie)	
3. Quarks (Gluonen) (Quantenchromodynamik)	

## Comparison: Atom --- Nucleus

Atom

$$H = \sum_{i=1}^Z \left( \frac{p_i^2}{2m} + \frac{Ze^2}{r_i} \right) + \sum_{i=1}^Z \sum_{j=i+1}^Z \frac{e^2}{|\vec{r}_i - \vec{r}_j|}$$

$\frac{Ze^2}{r_i}$  WW mit Kern  
 $\frac{e^2}{|\vec{r}_i - \vec{r}_j|}$  Elektri-Elektv.-WW

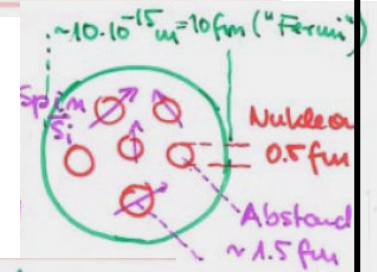


- Problem vollständig spezifiziert  
Coulomb WW bekannt  
(+ rel. Effekte; Quantenelektrodynamik;  $\ell$  + magn. Momente des Kerns, ...)
- Zentralfeld dominiert, Elektri-Elektv.-WW Schwach  
(in gewissem Sinne, + mittleres Feld der Elektronen)  
→ Periodensystem, Schalenmod.  $\sim 10^{-15} \text{ m} = 10 \text{ fm}$  ("Fermi")

Kern

$$H = \sum_{i=1}^A \frac{p_i^2}{2m} + \sum_{i=1}^A \sum_{j=i+1}^A V(\vec{r}_i - \vec{r}_j, \vec{s}_i, \vec{s}_j)$$

Nukl.-Nukl.-WW

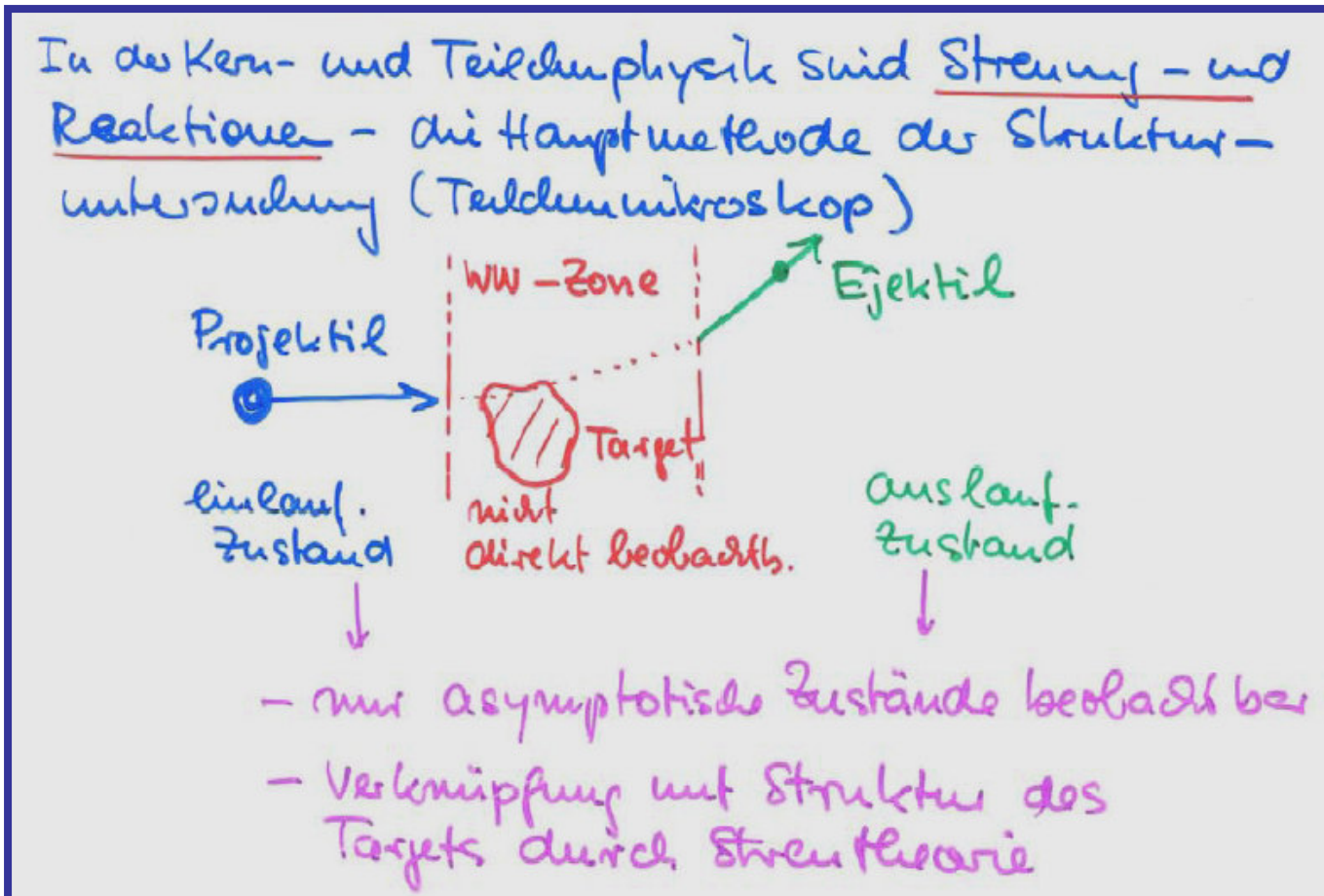


- $V_{NN}$  nicht exakt bekannt
- $V_{NN}$  nicht fundamental!
- kein dominierendes Zentralfeld,  
 $V_{NN}$  stark!  
dennoch: Schalenmodell möglich  
(Haxel, Jensen, Peopert-Flayer, 1949)
- Wichtigkeit von Modellen (Charakt. Moden des Systems, z. B. Phononen)  
↑  
Beziehung zu Vielteilchenkern

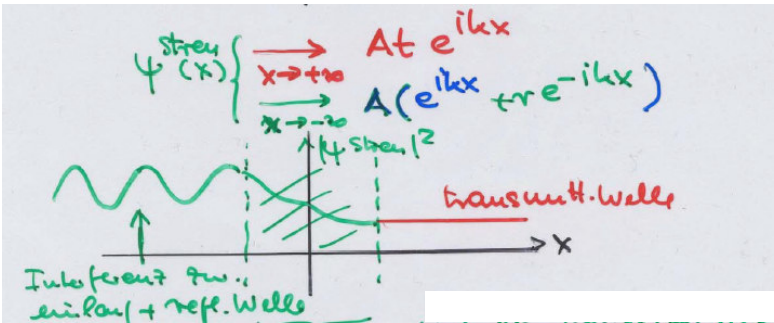
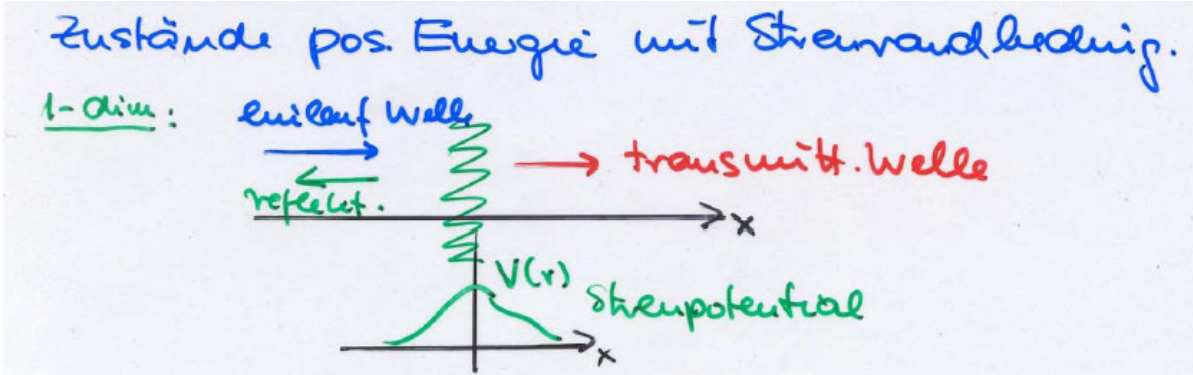
# 1. Reaction Theory

discuss essential aspects of theoretical treatment of scattering and reactions,

(this is general and not necessarily confined to nuclear reactions)



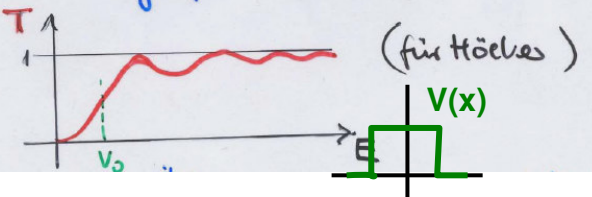
# Scattering states, Cross sections



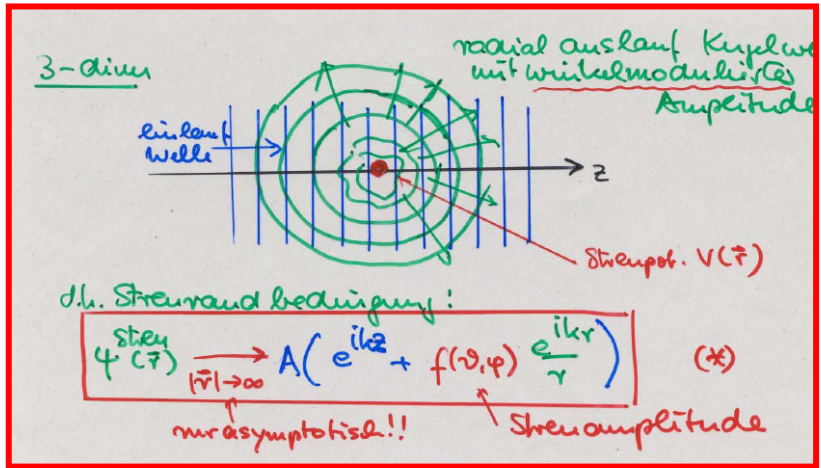
Stromdichten:  $\vec{j}_{\text{einf}} = A e^{i\vec{k}\vec{r}}$   
 $\vec{j} = \frac{\hbar}{m} \text{Im}(\psi^* \nabla \psi) = \frac{\hbar k}{m} |A|^2 = \vec{j}_{\text{einf}}$

$\vec{j}_{\text{trans}} = |t|^2 \frac{\hbar k}{m} |A|^2$ ;  $\vec{j}_{\text{refl}} = |r|^2 \frac{\hbar k}{m} |A|^2$

$T = \frac{\vec{j}_{\text{trans}}}{\vec{j}_{\text{einf}}} = |t|^2$  Transmiss. Koeff  
 $R = \frac{\vec{j}_{\text{refl}}}{\vec{j}_{\text{einf}}} = |r|^2$  Reflexionskoeff

$T$   (für Höcker)

$V(x)$



# Definition of cross section:

Proportionalität zwischen  $J_\Omega$  und  $J_0$ :

differentielle WQ  $\frac{d\sigma}{d\Omega} = \frac{J_\Omega}{J_0}$  gestreut. Teilchen in Raumwinkel  $d\Omega$  um  $\Omega$  und Zeit  
 einfall. Teilchen pro Fläche und Zeit

Totales Wirkungsquerschnitt:  $\sigma_{\text{tot}} = \int d\Omega \frac{d\sigma}{d\Omega}$

entspr. Verallg. für Reaktionen

a) klassisches WQ (rot. symm. Potentiale)

Teilchenzahl erhält:  $J_0 2\pi b db = J_\Omega d\Omega = J_0 2\pi \sin\theta d\theta$

$\frac{d\sigma^{\text{class}}}{d\Omega} = \frac{J_\Omega}{J_0} = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$

Zur Berechnung des klass. WQ muß die Beziehung  $\theta(b)$  oder  $b(\theta)$  bekannt sein:  
 Ableitungsfunktion:

$\sigma_{\text{tot}} = \int \frac{d\sigma}{d\Omega} d\Omega = \int \frac{b}{\sin\theta} \frac{db}{d\theta} 2\pi \sin\theta d\theta$   
 $= 2\pi \int_0^{b_{\text{max}}} b db = \begin{cases} \pi b_{\text{max}}^2 & \leftarrow \text{Fläche!} \\ \infty, \text{ falls } \theta \neq 0 \text{ für } b \rightarrow \infty \end{cases}$

b) Quantenmechanisches WQ

$\psi \rightarrow A(e^{ikz} + f(\Omega) \frac{e^{ikr}}{r})$

$J_0 = j^{\text{auf}} = |A|^2 \frac{\hbar k}{m}$   
 $J_\Omega = j^{\text{stre}} \frac{r^2 d\Omega}{d\Omega} = r^2 j^{\text{stre}} = |A|^2 |f|^2 \frac{\hbar k}{m} + O(\frac{1}{r})$

$\rightarrow \frac{d\sigma}{d\Omega} = |f(\Omega)|^2$  qu. diff. WQ

- Berechnung mit Wellenpaketen (WP):  
 $\rightarrow$  gleiches Ergebnis für realistische Parameter des WP

Damit ist WQ noch nicht berechnet! Dazu muß  $f(\Omega)$  berechnet werden:

Prinzip: SG mit Streurand bed.(x) lösen

- aus asympt. Verhalten  $f(\theta)$  extrahieren
- in 3-dim nicht einfach.



# Born Approximation

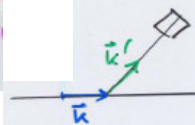
SG:  $(-\frac{\hbar^2}{2m} \Delta + V(r) - E) \psi(\vec{r}) = 0$   
 $(\Delta - u + k^2) \psi = 0$ ,  $E = \frac{\hbar^2 k^2}{2m}$ ,  $u = \frac{2m}{\hbar^2} V$   
 $(\Delta + k^2) \psi(\vec{r}) = U(r) \psi(\vec{r}) =: \text{Sinhom}(\vec{r})$

→ formal als inhom. Dgl. schreiben  
 → lösen mit Methode der Greenfunktion.

$$G^{(\pm)}(\vec{r}-\vec{r}') = -\frac{1}{4\pi} \frac{e^{\pm i k |\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|}$$

$$G^{(\pm)}(\vec{r}-\vec{r}') \rightarrow -\frac{1}{4\pi} \frac{e^{\pm i k r}}{r} e^{\mp i k \vec{r} \cdot \vec{r}'}$$

nichtige BB: Auslauf Kugel  
 also  $G^{(+)}$



Näherungslsg. der LS-Gleichung: Iteration

$$\begin{aligned} \psi^{(+)(0)} &= \psi_0(\vec{r}) && \mathcal{O}(V^0) \\ \psi^{(+)(1)} &= \psi_0(\vec{r}) + \int d^3r' G^{(+)}(\vec{r}-\vec{r}') U(\vec{r}') \psi^{(+)(0)}(\vec{r}') && \mathcal{O}(V^1) \\ \psi^{(+)(2)} &= \dots && \mathcal{O}(V^2) \end{aligned}$$

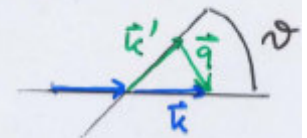
Born'sche Näherung: Verwende  $\psi^{(+)(0)}$  in Streuamplitude

$$f_k^{(BA)}(\Omega) = -\frac{1}{4\pi} \int d^3r e^{-i\vec{k}' \cdot \vec{r}} \left( \frac{2m}{\hbar^2} V(\vec{r}) \right) e^{i\vec{k} \cdot \vec{r}}$$

$$f_k^{(BA)}(\Omega) = -\frac{m}{2\pi\hbar^2} \int d^3r e^{i\vec{q} \cdot \vec{r}} V(\vec{r})$$

Born-Approx.

-  $\vec{q} = \vec{k} - \vec{k}'$  Impulsübertrag  
 $|\vec{q}| = q = 2k \sin \frac{\theta}{2}$   
 abhängig von Energie ( $k$ ) und Winkel.



Dann Lsg. der inhom. SG (S. 21)

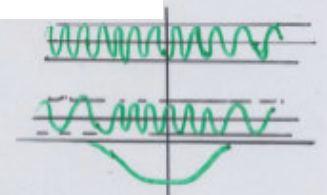
$$\psi^{(+)}(\vec{r}) = \underbrace{\psi_0(\vec{r})}_{e^{i\vec{k} \cdot \vec{r}}} + \int d^3r' G^{(+)}(\vec{r}-\vec{r}') U(\vec{r}') \underbrace{\psi^{(+)}(\vec{r}')}_{\text{Sinhom}}$$

Streuamplitude:

$$f_k(\Omega) = -\frac{1}{4\pi} \int d^3r e^{-i\vec{k}' \cdot \vec{r}} U(\vec{r}) \psi^{(+)}(\vec{r})$$

formale Lsg.!

- Gültigkeit der Näherung:  $V$  klein;  
 genau:  $V$  klein gegen  $E$ ,  
 Hochenergienäherung

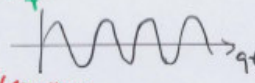


- Eigenschaft der FT: Breite von  $f$  und von  $F.T.$  umgekehrt prop

→ Streuampl. groß für  $q \lesssim \frac{1}{R}$   
 $2k \sin \frac{\theta}{2} \lesssim \frac{1}{R}$  }  $\left[ \theta \lesssim \frac{1}{kR} \right]$  Vorwärtswinkel  
 mit  $E \uparrow$

# Examples and Applications:

1) Streuung an Punktladung  $V_c = \frac{Ze^2}{r}$

$f_c^{(BA)}(q) = -\frac{2uZe^2}{k^2q} \int r dr \sin(qr) \frac{1}{r}$  

oszill., nicht konverg.

→ Konvergenzfaktor  $V \rightarrow V' = V e^{-\lambda r}, \lambda \rightarrow \infty$

$\sin(qr) = \frac{1}{2i}(e^{iqr} - e^{-iqr})$

→  $f_c^{(BA)}(q) = -\frac{2uZe^2}{k^2} \frac{1}{q^2} = f_{Rukh}(q)$  Rutherford-Amplitude

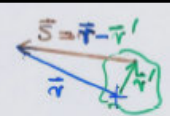
$\frac{d\sigma^{(BA)}}{d\Omega} = |f_c^{(BA)}|^2 = \frac{4u^2(Ze^2)^2}{(k^2q^2)^2} = \left(\frac{2e^2}{4E}\right)^2 \frac{1}{\sin^4 \frac{\theta}{2}} = \frac{d\sigma}{d\Omega} |Rukh|$

→ für Coulomb:  $\frac{d\sigma}{d\Omega}^{class} \equiv \frac{d\sigma}{d\Omega}^{(BA)} = \frac{d\sigma}{d\Omega}^{Rukh}$  ! (Zufall)

2) endlich Ladungsverteil.

$V(\vec{r}) = \int d^3r' \rho(\vec{r}') V_c(\vec{r}-\vec{r}')$

$f_V^{(BA)}(\Omega) = -\frac{u}{2\pi k^2} \int d^3r' e^{i\vec{q}\cdot\vec{r}'} V(\vec{r})$

$\vec{s} = \vec{r} - \vec{r}'$    $\rho(\vec{r}')$

$\vec{s} + \vec{r}' = \vec{r} = \int d^3r' \rho(\vec{r}') V_c(\vec{s})$

$= \int d^3r' e^{i\vec{q}\cdot\vec{r}'} \rho(\vec{r}') \int d^3s e^{i\vec{q}\cdot\vec{s}} V_c(\vec{s}) \left(-\frac{u}{2\pi k^2}\right)$

$=: F(\vec{q})$   $\uparrow$   $f_{Rukh}!$

$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega} |Rukh| |F(\vec{q})|^2$

mit  $F(\vec{q}) = \int d^3r' e^{i\vec{q}\cdot\vec{r}'} \rho(\vec{r}')$  Form faktor  
 $\hat{=}$  Fouriertransf der Ladungsverteil.

→ Messung der Abt. ver!

4) Bestimmung der Ladungsverteilung

$\alpha)$  Inverse Fouriertransformation

$\rho(\vec{r}) = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} e^{-i\vec{q}\cdot\vec{r}} F(\vec{q}) d^3q$

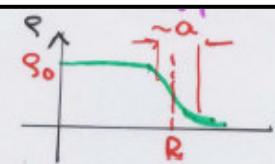
Belast  $\vec{q}$ ? (d.h. Winkel, Energie)

→ Fehlerband (Fourier-Bessel Analyse)

$\beta)$  Ansatz für  $\rho(\vec{r})$  und Bestimmung der Parameter

i) Nukleon  $\rho_c = \frac{\rho_0}{8\pi} e^{-ar} \rightarrow F(q^2) = \frac{1}{(1+q^2/a^2)^2}$

experiment  $a \approx 4.2 \text{ fm} \rightarrow \sqrt{\langle r^2 \rangle} \approx 0.86 \text{ fm}$

ii) Kern  $\rho_c(r) = \frac{\rho_0}{1 + \exp\left(\frac{r-R}{a}\right)}$  

exp.:

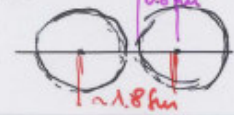
- $R = r_0 A^{1/3}, r_0 \approx \text{const} \approx 1.15 \text{ fm}$
- $a \approx \text{const} \approx 0.6 \text{ fm}$
- $\rho_0 \approx \text{const} \approx 0.16 \text{ fm}^{-3} \approx \frac{1}{6} \frac{\text{Nukl}}{\text{fm}^3}$

Flüssigkeitseigenschaften, Sättigung!

$V \propto R^3 \sim A, \rho \ll R$  (außer für kleine Kerne)

$\frac{Vol}{Nukl} = \frac{V}{A} = \frac{4\pi}{3} r_0^3 \rightarrow$  mittl. Abstand der Nukl.

$\bar{d} \approx \sqrt[3]{\frac{V}{A}} \sim 1.6 \text{ fm} \sim 1.8 \text{ fm}$

 → dichte Packung!

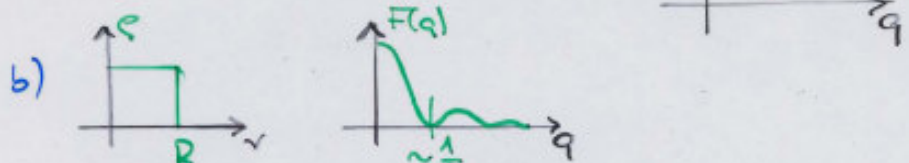
# Applications of BA (contd.)

5) Qualitatives Verhalten:

$$a) F(\vec{q}) = \int d^3r \rho(\vec{r}) e^{i\vec{q}\cdot\vec{r}}$$

$$= 1 - \frac{q^2}{6} \langle r^2 \rangle + \dots$$

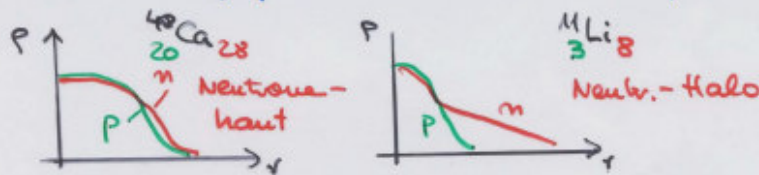
$$1 + i\vec{q}\cdot\vec{r} - \frac{1}{2}(\vec{q}\cdot\vec{r})^2 + \dots$$



6) Elektronenstrom, mit Ladungsverteilung, Scherze Massenverteilung?

Bei stabilen Kernen approx:  $\rho_{mass} = \frac{A}{Z}$  Scherze

aber interessante Effekte nur bei Kernen mit großen Neutronenüberschuß (exot. Ker)



Messung  $\rho_{ch}$ : Hadronstreuung: p,  $\pi$ ,  $\rightarrow$  "optisches" Pot.  
 Schwache WW:  $\nu \leftarrow$  (Vorsicht!)

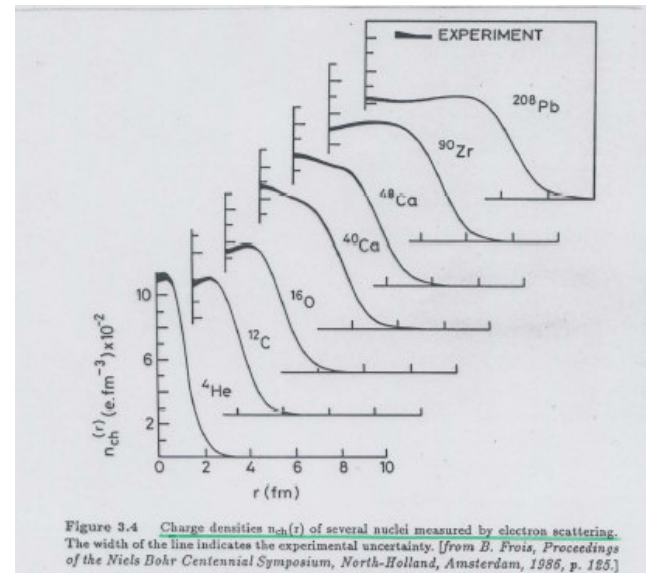
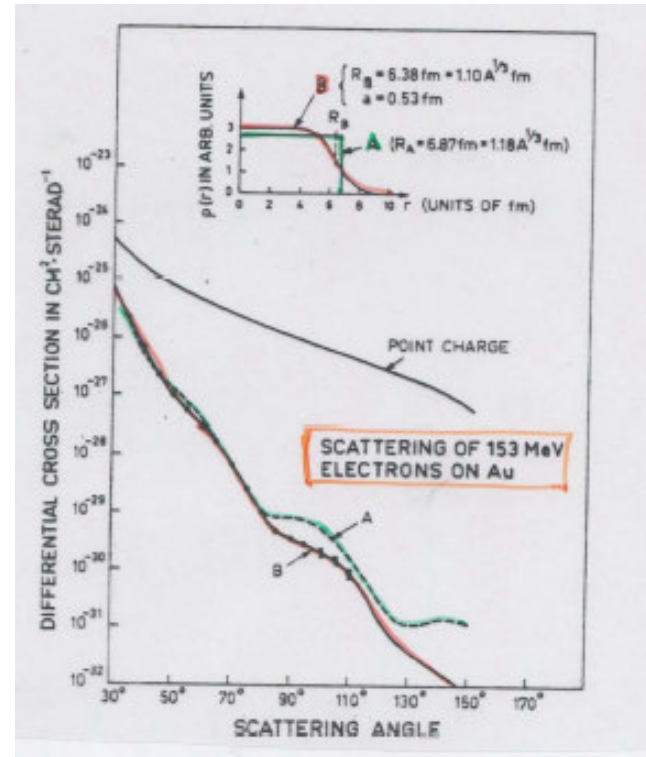


Figure 3.4 Charge densities  $n_{ch}(r)$  of several nuclei measured by electron scattering. The width of the line indicates the experimental uncertainty. [from B. Frois, Proceedings of the Niels Bohr Centennial Symposium, North-Holland, Amsterdam, 1986, p. 125.]